

# Unipotent Brauer character values of $GL(n, \mathbb{F}_q)$ and the forgotten basis of the Hall algebra <sup>\*</sup>

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## Abstract

We give a formula for the values of irreducible unipotent  $p$ -modular Brauer characters of  $GL(n, \mathbb{F}_q)$  at unipotent elements, where  $p$  is a prime not dividing  $q$ , in terms of (unknown!) weight multiplicities of quantum  $GL_n$  and certain generic polynomials  $S_{\lambda, \mu}(q)$ . These polynomials arise as entries of the transition matrix between the renormalized Hall-Littlewood symmetric functions and the forgotten symmetric functions. We also provide an alternative combinatorial algorithm working in the Hall algebra for computing  $S_{\lambda, \mu}(q)$ .

## 1 Introduction

In the character theory of the finite general linear group  $G_n = GL(n, \mathbb{F}_q)$ , the *Gelfand-Graev character*  $\Gamma_n$  plays a fundamental role. By definition [5],  $\Gamma_n$  is the character obtained by inducing a “general position” linear character from a maximal unipotent subgroup. It has support in the set of unipotent elements of  $G_n$  and for a unipotent element  $u$  of type  $\lambda$  (i.e. the block sizes of the Jordan normal form of  $u$  are the parts of the partition  $\lambda$ ) Kawanaka [7, 3.2.24] has shown that

$$\Gamma_n(u) = (-1)^n(1-q)(1-q^2)\dots(1-q^{h(\lambda)}), \quad (1.1)$$

where  $h(\lambda)$  is the number of non-zero parts of  $\lambda$ . The starting point for this article is the problem of calculating the operator determined by Harish-Chandra multiplication by  $\Gamma_n$ .

We have restricted our attention throughout to character values at unipotent elements, when it is convenient to work in terms of the *Hall algebra*, that is [13, §10.1], the vector space  $\mathfrak{g} = \bigoplus_{n \geq 0} \mathfrak{g}_n$ , where  $\mathfrak{g}_n$  denotes the set of unipotent-supported  $\mathbb{C}$ -valued class functions on  $G_n$ , with multiplication coming from the Harish-Chandra induction operator. For a partition  $\lambda$  of  $n$ , let  $\pi_\lambda \in \mathfrak{g}_n$  denote the class function which is 1 on unipotent elements of type  $\lambda$  and zero on all other conjugacy classes of  $G_n$ . Then,  $\{\pi_\lambda\}$  is a basis for the Hall algebra labelled by all partitions. Let  $\gamma_n : \mathfrak{g} \rightarrow \mathfrak{g}$  be the linear operator determined by multiplication in  $\mathfrak{g}$  by  $\Gamma_n$ . We describe in §2 an explicit recursive algorithm, involving the combinatorics of addable and removable nodes, for calculating the effect of  $\gamma_n$  on the basis  $\{\pi_\lambda\}$ . As an illustration of the algorithm, we rederive Kawanaka’s formula (1.1) in Example 2.12.

Now recall from [13] that  $\mathfrak{g}$  is isomorphic to the algebra  $\Lambda_{\mathbb{C}}$  of symmetric functions over  $\mathbb{C}$ , the isomorphism sending the basis element  $\pi_\lambda$  of  $\mathfrak{g}$  to the Hall-Littlewood symmetric function  $\tilde{P}_\lambda \in \Lambda_{\mathbb{C}}$  (renormalized as in [9, §II.3, ex.2]). Consider instead the

<sup>\*</sup>1991 subject classification: 20C20, 05E05, 20C33.

<sup>‡</sup>Partially supported by NSF grant no. DMS-9801442.

element  $\vartheta_\lambda \in \mathfrak{g}$  which maps under this isomorphism to the *forgotten symmetric function*  $f_\lambda \in \Lambda_{\mathbb{C}}$  (see [9, §I.2]). Introduce the renormalized Gelfand-Graev operator  $\hat{\gamma}_n = \delta \circ \gamma_n$ , where  $\delta : \mathfrak{g} \rightarrow \mathfrak{g}$  is the linear map with  $\delta(\pi_\lambda) = \frac{1}{q^{h(\lambda)-1}} \pi_\lambda$  for all partitions  $\lambda$ . We show in Theorem 3.5 that

$$\vartheta_\lambda = \sum_{(n_1, \dots, n_h)} \hat{\gamma}_{n_1} \circ \hat{\gamma}_{n_2} \circ \dots \circ \hat{\gamma}_{n_h}(\pi_{(0)}), \quad (1.2)$$

summing over all  $(n_1, \dots, n_h)$  obtained by reordering the non-zero parts  $\lambda_1, \dots, \lambda_h$  of  $\lambda$  in all possible ways. Thus, we obtain a direct combinatorial construction of the ‘forgotten basis’  $\{\vartheta_\lambda\}$  of the Hall algebra.

Let  $K = (K_{\lambda, \mu})$  denote the matrix of Kostka numbers [9, I, (6.4)],  $\tilde{K} = (\tilde{K}_{\lambda, \mu}(q))$  denote the matrix of Kostka-Foulkes polynomials (renormalized as in [9, III, (7.11)]) and  $J = (J_{\lambda, \mu})$  denote the matrix with  $J_{\lambda, \mu} = 0$  unless  $\mu = \lambda'$  when it is 1, where  $\lambda'$  is the conjugate partition to  $\lambda$ . Consulting [9, §I.6, §III.6], the transition matrix between the bases  $\{\pi_\lambda\}$  and  $\{\vartheta_\lambda\}$ , i.e. the matrix  $S = (S_{\lambda, \mu}(q))$  of coefficients such that

$$\vartheta_\lambda = \sum_{\mu} S_{\lambda, \mu}(q) \pi_\mu, \quad (1.3)$$

is then given by the formula  $S = K^{-1} J \tilde{K}$ ; in particular, this implies that  $S_{\lambda, \mu}(q)$  is a polynomial in  $q$  with integer coefficients. Our alternative approach to computing  $\vartheta_\lambda$  using (1.2) allows explicit computation of the polynomials  $S_{\lambda, \mu}(q)$  in some extra cases (e.g. when  $\mu = (1^n)$ ) not easily deduced from the matrix product  $K^{-1} J \tilde{K}$ .

To explain our interest in this, let  $\chi_\lambda$  denote the irreducible unipotent character of  $G_n$  labelled by the partition  $\lambda$ , as constructed originally in [12], and let  $\sigma_\lambda \in \mathfrak{g}$  denote its projection to unipotent-supported class functions. So,  $\sigma_\lambda$  is the element of  $\mathfrak{g}$  mapping to the Schur function  $s_\lambda$  under the isomorphism  $\mathfrak{g} \rightarrow \Lambda_{\mathbb{C}}$  (see [13]). Since  $\sigma_{\lambda'} = \sum_{\mu} K_{\lambda, \mu} \vartheta_\mu$  [9, §I.6], we deduce that the value of  $\chi_{\lambda'}$  at a unipotent element  $u$  of type  $\nu$  can be expressed in terms of the Kostka numbers  $K_{\lambda, \mu}$  and the polynomials  $S_{\mu, \nu}(q)$  as

$$\chi_{\lambda'}(u) = \sum_{\mu} K_{\lambda, \mu} S_{\mu, \nu}(q). \quad (1.4)$$

This is a rather clumsy way of expressing the unipotent character values in the ordinary case, but this point of view turns out to be well-suited to describing the irreducible unipotent *Brauer characters*.

So now suppose that  $p$  is a prime not dividing  $q$ ,  $\mathbb{k}$  is a field of characteristic  $p$  and let the multiplicative order of  $q$  modulo  $p$  be  $\ell$ . In [6], James constructed for each partition  $\lambda$  of  $n$  an absolutely irreducible, unipotent  $\mathbb{k}G_n$ -module  $D_\lambda$  (denoted  $L(1, \lambda)$  in [1]), and showed that the set of all  $D_\lambda$  gives the complete set of non-isomorphic irreducible modules that arise as constituents of the permutation representation of  $\mathbb{k}G_n$  on cosets of a Borel subgroup. Let  $\chi_\lambda^p$  denote the Brauer character of the module  $D_\lambda$ , and  $\sigma_\lambda^p \in \mathfrak{g}$  denote the projection of  $\chi_\lambda^p$  to unipotent-supported class functions. Then, as a direct consequence of the results of Dipper and James [3], we show in Theorem 4.6

that  $\sigma_{\lambda'}^p = \sum_{\mu} K_{\lambda,\mu}^{p,\ell} \vartheta_{\mu}$  where  $K_{\lambda,\mu}^{p,\ell}$  denotes the weight multiplicity of the  $\mu$ -weight space in the irreducible high-weight module of high-weight  $\lambda$  for *quantum*  $GL_n$ , at an  $\ell$ th root of unity over a field of characteristic  $p$ . In other words, for a unipotent element  $u$  of type  $\nu$ , we have the modular analogue of (1.4):

$$\chi_{\lambda'}^p(u) = \sum_{\mu} K_{\lambda,\mu}^{p,\ell} S_{\mu,\nu}(q) \quad (1.5)$$

This formula reduces the problem of calculating the values of the irreducible unipotent Brauer characters at unipotent elements to knowing the modular Kostka numbers  $K_{\lambda,\mu}^{p,\ell}$  and the polynomials  $S_{\mu,\nu}(q)$ .

Most importantly, taking  $\nu = (1^n)$  in (1.5), we obtain the degree formula:

$$\chi_{\lambda'}^p(1) = \sum_{\mu} K_{\lambda,\mu}^{p,\ell} S_{\mu,(1^n)}(q) \quad (1.6)$$

where, as a consequence of (1.2) (see Example 3.7),

$$S_{\mu,(1^n)}(q) = \sum_{(n_1, \dots, n_h)} \left[ \prod_{i=1}^n (q^i - 1) / \prod_{i=1}^h (q^{n_1 + \dots + n_i} - 1) \right] \quad (1.7)$$

summing over all  $(n_1, \dots, n_h)$  obtained by reordering the non-zero parts  $\mu_1, \dots, \mu_h$  of  $\mu$  in all possible ways. This formula was first proved in [1, §5.5], as a consequence of a result which can be regarded as the modular analogue of Zelevinsky's branching rule [13, §13.5] involving the affine general linear group. The proof presented here is independent of [1] (excepting some self-contained results from [1, §5.1]), appealing instead directly to the original characteristic 0 branching rule of Zelevinsky, together with the work of Dipper and James on decomposition matrices. We remark that since all of the integers  $S_{\mu,(1^n)}(q)$  are positive, the formula (1.6) can be used to give quite powerful *lower bounds* for the degrees of the irreducible Brauer characters, by exploiting a  $q$ -analogue of the Premet-Suprunenko bound for the  $K_{\lambda,\mu}^{p,\ell}$ . The details can be found in [2].

To conclude this introduction, we list in the table below the polynomials  $S_{\lambda,\mu}(q)$  for  $n \leq 4$ :

1	1	2	1 <sup>2</sup>	3	21	1 <sup>3</sup>
1	1	2	-1    q-1	3	1    1-q    (q <sup>2</sup> -1)(q-1)	(q <sup>2</sup> -1)(q-1)
1 <sup>2</sup>	1	1 <sup>2</sup>	1	21	-2    q-2    (q-1)(q+2)	(q-1)(q+2)
1 <sup>3</sup>	1	1 <sup>3</sup>	1	1 <sup>3</sup>	1	1
4	31	2 <sup>2</sup>	21 <sup>2</sup>	4	31	1 <sup>4</sup>
4	-1    q-1	q-1	(1-q)(q <sup>2</sup> -1)	(q <sup>3</sup> -1)(q <sup>2</sup> -1)(q-1)	(q <sup>3</sup> -1)(q <sup>2</sup> -1)(q-1)	(q <sup>3</sup> -1)(q <sup>2</sup> -1)(q-1)
31	2    2-q	(1-q)(q+2)	(q <sup>2</sup> -1)(q-2)	(q <sup>2</sup> -1)(q <sup>3</sup> +q <sup>2</sup> -2)	(q <sup>2</sup> -1)(q <sup>3</sup> +q <sup>2</sup> -2)	(q <sup>2</sup> -1)(q <sup>3</sup> +q <sup>2</sup> -2)
2 <sup>2</sup>	1    1-q	q <sup>2</sup> -q+1	1-q	(q <sup>3</sup> -1)(q-1)	(q <sup>3</sup> -1)(q-1)	(q <sup>3</sup> -1)(q-1)
21 <sup>2</sup>	-3    q-3	q-3	q <sup>2</sup> +q-3	q <sup>3</sup> +q <sup>2</sup> +q-3	q <sup>3</sup> +q <sup>2</sup> +q-3	q <sup>3</sup> +q <sup>2</sup> +q-3
1 <sup>4</sup>	1	1	1	1	1	1

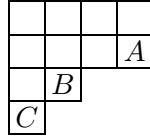
Finally, I would like to thank Alexander Kleshchev for several helpful discussions about this work.

## 2 An algorithm for computing $\gamma_n$

We will write  $\lambda \vdash n$  to indicate that  $\lambda$  is a partition of  $n$ , that is, a sequence  $\lambda = (\lambda_1 \geq \lambda_2 \geq \dots)$  of non-negative integers summing to  $n$ . Given  $\lambda \vdash n$ , we denote its *Young diagram* by  $[\lambda]$ ; this is the set of *nodes*

$$[\lambda] = \{(i, j) \in \mathbb{N} \times \mathbb{N} \mid 1 \leq j \leq \lambda_i\}.$$

By an *addable node* (for  $\lambda$ ), we mean a node  $A \in \mathbb{N} \times \mathbb{N}$  such that  $[\lambda] \cup \{A\}$  is the diagram of a partition; we denote the new partition obtained by adding the node  $A$  to  $\lambda$  by  $\lambda \cup A$ . By a *removable node* (for  $\lambda$ ) we mean a node  $B \in [\lambda]$  such that  $[\lambda] \setminus \{B\}$  is the diagram of a partition; we denote the new partition obtained by removing  $B$  from  $\lambda$  by  $\lambda \setminus B$ . The *depth*  $d(B)$  of the node  $B = (i, j) \in \mathbb{N} \times \mathbb{N}$  is the row number  $i$ . If  $B$  is removable for  $\lambda$ , it will also be convenient to define  $e(B)$  (depending also on  $\lambda$ !) to be the depth of the next removable node above  $B$  in the partition  $\lambda$ , or 0 if no such node exists. For example consider the partition  $\lambda = (4, 4, 2, 1)$ , and let  $A, B, C$  be the removable nodes in order of increasing depth:



Then,  $e(A) = 0, e(B) = d(A) = 2, e(C) = d(B) = 3, d(C) = 4$ .

Now fix a prime power  $q$  and let  $G_n$  denote the finite general linear group  $GL(n, \mathbb{F}_q)$  as in the introduction. Let  $V_n = \mathbb{F}_q^n$  denote the natural  $n$ -dimensional left  $G_n$ -module, with standard basis  $v_1, \dots, v_n$ . Let  $H_n$  denote the *affine general linear group*  $AGL(n, \mathbb{F}_q)$ . This is the semidirect product  $H_n = V_n G_n$  of  $G_n$  acting on the elementary Abelian group  $V_n$ . We always work with the standard embedding  $H_n \hookrightarrow G_{n+1}$  that identifies  $H_n$  with the subgroup of  $G_{n+1}$  consisting of all matrices of the form:

$$\left[ \begin{array}{ccc|c} & & & \\ & * & & * \\ & & & \\ \hline 0 & \dots & 0 & 1 \end{array} \right].$$

Thus, we have a chain of subgroups  $1 = H_0 \subset G_1 \subset H_1 \subset G_2 \subset H_2 \subset \dots$  (by convention, we allow the notations  $G_0, H_0$  and  $V_0$ , all of which denote groups with one element.)

For  $\lambda \vdash n$ , let  $u_\lambda \in G_n \subset H_n$  denote the upper uni-triangular matrix consisting of Jordan blocks of sizes  $\lambda_1, \lambda_2, \dots$  down the diagonal. As is well-known,  $\{u_\lambda \mid \lambda \vdash n\}$  is a set of representatives of the unipotent conjugacy classes in  $G_n$ . We wish to describe instead the unipotent classes in  $H_n$ . These were determined in [10, §1], but the notation here will be somewhat different. For  $\lambda \vdash n$  and an addable node  $A$  for  $\lambda$ , define the upper uni-triangular  $(n+1) \times (n+1)$  matrix  $u_{\lambda, A} \in H_n \subset G_{n+1}$  by

$$u_{\lambda, A} = \begin{cases} u_\lambda & \text{if } A \text{ is the deepest addable node,} \\ v_{\lambda_1 + \dots + \lambda_{d(A)}} u_\lambda & \text{otherwise.} \end{cases}$$

If instead  $\lambda \vdash (n+1)$  and  $B$  is removable for  $\lambda$  (hence addable for  $\lambda \setminus B$ ) define  $u_{\lambda,B}$  to be a shorthand for  $u_{\lambda \setminus B, B} \in H_n$ . To aid translation between our notation and that of [10], we note that  $u_{\lambda,B}$  is conjugate to the element denoted  $c_{n+1}(1^{(k)}, \mu)$  there, where  $k = \lambda_{d(B)}$  and  $\mu$  is the partition obtained from  $\lambda$  by removing the  $d(B)$ th row. Then:

**2.1. Lemma.** (i) *The set*

$$\{u_{\lambda,A} \mid \lambda \vdash n, A \text{ addable for } \lambda\} = \{u_{\lambda,B} \mid \lambda \vdash (n+1), B \text{ removable for } \lambda\}$$

*is a set of representatives of the unipotent conjugacy classes of  $H_n$ .*

(ii) *For  $\lambda \vdash (n+1)$  and a removable node  $B$ ,  $|C_{G_{n+1}}(u_{\lambda})|/|C_{H_n}(u_{\lambda,B})| = q^{d(B)} - q^{e(B)}$ .*

*Proof.* Part (i) is a special case of [10, 1.3(i)], where all conjugacy classes of the group  $H_n$  are described. For (ii), combine the formula for  $|C_{G_{n+1}}(u_{\lambda})|$  from [12, 2.2] with [10, 1.3(ii)], or calculate directly.  $\square$

For any group  $G$ , we write  $C(G)$  for the set of  $\mathbb{C}$ -valued class functions on  $G$ . Let  $\mathfrak{g} = \bigoplus_{n \geq 0} \mathfrak{g}_n \subset \bigoplus_{n \geq 0} C(G_n)$  denote the Hall algebra as in the introduction. We recall that  $\mathfrak{g}$  is a graded Hopf algebra in the sense of [13], with multiplication  $\diamond : \mathfrak{g} \otimes \mathfrak{g} \rightarrow \mathfrak{g}$  arising from Harish-Chandra induction and comultiplication  $\Delta : \mathfrak{g} \rightarrow \mathfrak{g} \otimes \mathfrak{g}$  arising from Harish-Chandra restriction, see [13, §10.1] for fuller details. Also defined in the introduction,  $\mathfrak{g}$  has the natural ‘characteristic function’ basis  $\{\pi_{\lambda}\}$  labelled by all partitions.

By analogy, we introduce an extended version of the Hall algebra corresponding to the affine general linear group. This is the vector space  $\mathfrak{h} = \bigoplus_{n \geq 0} \mathfrak{h}_n$  where  $\mathfrak{h}_n$  is the subspace of  $C(H_n)$  consisting of all class functions with support in the set of unipotent elements of  $H_n$ , with algebra structure to be explained below. To describe a basis for  $\mathfrak{h}$ , given  $\lambda \vdash n$  and an addable node  $A$ , define  $\pi_{\lambda,A} \in C(H_n)$  to be the class function which takes value 1 on  $u_{\lambda,A}$  and is zero on all other conjugacy classes of  $H_n$ . Given  $\lambda \vdash (n+1)$  and a removable  $B$ , set  $\pi_{\lambda,B} = \pi_{\lambda \setminus B, B}$ . Then, in view of Lemma 2.1(i),  $\{\pi_{\lambda,A} \mid \lambda \text{ a partition, } A \text{ addable for } \lambda\} = \{\pi_{\lambda,B} \mid \lambda \text{ a partition, } B \text{ removable for } \lambda\}$  is a basis for  $\mathfrak{h}$ .

Now we introduce various operators as in [13, §13.1] (but take notation instead from [1, §5.1]). First, for  $n \geq 0$ , we have the inflation operator

$$e_0^n : C(G_n) \rightarrow C(H_n)$$

defined by  $(e_0^n \chi)(vg) = \chi(g)$  for  $\chi \in C(G_n), v \in V_n, g \in G_n$ . Next, fix a non-trivial additive character  $\chi : \mathbb{F}_q \rightarrow \mathbb{C}^\times$  and let  $\chi_n : V_n \rightarrow \mathbb{C}^\times$  be the character defined by  $\chi_n(\sum_{i=1}^n c_i v_i) = \chi(c_n)$ . The group  $G_n$  acts naturally on the characters  $C(V_n)$  and one easily checks that the subgroup  $H_{n-1} < G_n$  centralizes  $\chi_n$ . In view of this, it makes sense to define for each  $n \geq 1$  the operator

$$e_+^n : C(H_{n-1}) \rightarrow C(H_n),$$

namely, the composite of inflation from  $H_{n-1}$  to  $V_n H_{n-1}$  with the action of  $V_n$  being via the character  $\chi_n$ , followed by ordinary induction from  $V_n H_{n-1}$  to  $H_n$ . Finally, for  $n \geq 1$  and  $1 \leq i \leq n$ , we have the operator

$$e_i^n : C(G_{n-i}) \rightarrow C(H_n)$$

defined inductively by  $e_i^n = e_+^n \circ e_{i-1}^{n-1}$ . The significance of these operators is due to the following lemma [13, §13.2]:

**2.2. Lemma.** *The operator  $e_0^n \oplus e_1^n \oplus \cdots \oplus e_n^n : C(G_n) \oplus C(G_{n-1}) \oplus \cdots \oplus C(G_0) \rightarrow C(H_n)$  is an isometry.*

We also have the usual restriction and induction operators

$$\text{res}_{H_{n-1}}^{G_n} : C(G_n) \rightarrow C(H_{n-1}) \quad \text{ind}_{H_{n-1}}^{G_n} : C(H_{n-1}) \rightarrow C(G_n).$$

One checks that all of the operators  $e_0^n, e_+^n, \text{res}_{H_{n-1}}^{G_n}$  and  $\text{ind}_{H_{n-1}}^{G_n}$  send class functions with unipotent support to class functions with unipotent support. So, we can define the following operators between  $\mathfrak{g}$  and  $\mathfrak{h}$ , by restricting the operators listed to unipotent-supported class functions:

$$e_+ : \mathfrak{h} \rightarrow \mathfrak{h}, \quad e_+ = \bigoplus_{n \geq 1} e_+^n; \quad (2.3)$$

$$e_i : \mathfrak{g} \rightarrow \mathfrak{h}, \quad e_i = \bigoplus_{n \geq i} e_i^n = (e_+)^i \circ e_0; \quad (2.4)$$

$$\text{ind} : \mathfrak{h} \rightarrow \mathfrak{g}, \quad \text{ind} = \bigoplus_{n \geq 1} \text{ind}_{H_{n-1}}^{G_n}; \quad (2.5)$$

$$\text{res} : \mathfrak{g} \rightarrow \mathfrak{h}, \quad \text{res} = \bigoplus_{n \geq 0} \text{res}_{H_{n-1}}^{G_n} \quad (2.6)$$

where for the last definition,  $\text{res}_{H_{-1}}^{G_0}$  should be interpreted as the zero map.

Now we indicate briefly how to make  $\mathfrak{h}$  into a graded Hopf algebra. In view of Lemma 2.2, there are unique linear maps  $\diamond : \mathfrak{h} \otimes \mathfrak{h} \rightarrow \mathfrak{h}$  and  $\Delta : \mathfrak{h} \rightarrow \mathfrak{h} \otimes \mathfrak{h}$  such that

$$(e_i \chi) \diamond (e_j \tau) = e_{i+j}(\chi \diamond \tau), \quad (2.7)$$

$$\Delta(e_k \psi) = \sum_{i+j=k} (e_i \otimes e_j) \Delta(\psi), \quad (2.8)$$

for all  $i, j, k \geq 0$  and  $\chi, \tau, \psi \in \mathfrak{g}$ . One can check, using Lemma 2.2 and the fact that  $\mathfrak{g}$  is a graded Hopf algebra, that these operations endow  $\mathfrak{h}$  with the structure of a graded Hopf algebra (the unit element is  $e_0 \pi_{(0)}$ , and the counit is the map  $e_i \chi \mapsto \delta_{i,0} \varepsilon(\chi)$  where  $\varepsilon : \mathfrak{g} \rightarrow \mathbb{C}$  is the counit of  $\mathfrak{g}$ ). Moreover, the map  $e_0 : \mathfrak{g} \rightarrow \mathfrak{h}$  is a Hopf algebra embedding. Unlike for the operations of  $\mathfrak{g}$ , we do not know of a natural representation theoretic interpretation for these operations on  $\mathfrak{h}$  except in special cases, see [1, §5.2].

The effect of the operators (2.3)–(2.6) on our characteristic function bases is described explicitly by the following lemma:

**2.9. Lemma.** *Let  $\lambda \vdash n$  and label the addable nodes (resp. removable nodes) of  $\lambda$  as  $A_1, A_2, \dots, A_s$  (resp.  $B_1, B_2, \dots, B_{s-1}$ ) in order of increasing depth. Also let  $B = B_r$  be some fixed removable node. Then,*

- (i)  $e_0\pi_\lambda = \sum_{i=1}^s \pi_{\lambda, A_i};$
- (ii)  $e_+\pi_{\lambda, B} = q^{d(B)} \sum_{i=r+1}^s \pi_{\lambda, A_i} - q^{e(B)} \sum_{i=r}^s \pi_{\lambda, A_i}.$
- (iii)  $\text{res } \pi_\lambda = \sum_{i=1}^{s-1} \pi_{\lambda, B_i};$
- (iv)  $\text{ind } \pi_{\lambda, B} = (q^{d(B)} - q^{e(B)})\pi_\lambda;$
- (v)  $\text{ind} \circ \text{res } \pi_\lambda = (q^{h(\lambda)} - 1)\pi_\lambda.$

*Proof.* (i) For  $\mu \vdash n$  and  $A$  addable, we have by definition that  $(e_0\pi_\lambda)(u_{\mu, A}) = \pi_\lambda(u_\mu) = \delta_{\lambda, \mu}$ . Hence,  $e_0\pi_\lambda = \sum_A \pi_{\lambda, A}$ , summing over all addable nodes  $A$  for  $\lambda$ .

(ii) This is a special case of [10, 2.4] translated into our notation.

(iii) For  $\mu \vdash n$  and  $B$  removable,  $(\text{res } \pi_\lambda)(u_{\mu, B})$  is zero unless  $u_{\mu, B}$  is conjugate in  $G_n$  to  $u_\lambda$ , when it is one. So the result follows on observing that  $u_{\mu, B}$  is conjugate in  $G_n$  to  $u_\lambda$  if and only if  $\mu = \lambda$ .

(iv) We can write  $\text{ind } \pi_{\lambda, B} = \sum_{\mu \vdash n} c_\mu \pi_\mu$ . To calculate the coefficient  $c_\mu$  for fixed  $\mu \vdash n$ , we use (iii), Lemma 2.1(ii) and Frobenius reciprocity.

(v) This follows at once from (iii) and (iv) since  $\sum_B (q^{d(B)} - q^{e(B)}) = q^{h(\lambda)} - 1$ , summing over all removable nodes  $B$  for  $\lambda$ .  $\square$

Lemma 2.9(i),(ii) give explicit formulae for computing the operator  $e_n = (e_+)^n \circ e_0$ . The connection between  $e_n$  and the Gelfand-Graev operator  $\gamma_n$  defined in the introduction comes from the following result:

**2.10. Theorem.** For  $n \geq 1$ ,  $\gamma_n = \text{ind} \circ e_{n-1}$ .

*Proof.* In [1, Theorem 5.1e], we showed directly from the definitions that for any  $\chi \in C(G_m)$  and any  $n \geq 1$ , the class function  $\chi \cdot \Gamma_n \in C(G_{m+n})$  obtained by Harish-Chandra induction from  $(\chi, \Gamma_n) \in C(G_m) \times C(G_n)$  is equal to  $\text{res}_{G_{m+n}}^{H_{m+n}}(e_n^{m+n}\chi)$ . Moreover, by [1, Lemma 5.1c(iii)], we have that  $\text{res}_{G_{m+n}}^{H_{m+n}} \circ e_+^{m+n} = \text{ind}_{H_{m+n-1}}^{G_{m+n}}$ . Hence,

$$\chi \cdot \Gamma_n = \text{ind}_{H_{m+n-1}}^{G_{m+n}} (e_{n-1}^{m+n-1}\chi).$$

The theorem is just a restatement of this formula at the level of unipotent-supported class functions.  $\square$

**2.11. Example.** We show how to calculate  $\gamma_2\pi_{(3,2)}$  using Lemma 2.9 and the theorem. We omit the label  $\pi$  in denoting basis elements, and in the case of the intermediate basis

elements of  $\mathfrak{h}$ , we mark removable nodes with  $\times$ .

$$\begin{aligned}
\gamma_2 \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \square & \square & \square \\ \hline \end{array} &= \text{ind } \circ e_+ \left( \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \times \\ \hline \square & \square & \square & \square \\ \hline \end{array} + \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \square & \square & \times \\ \hline \end{array} + \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \square & \square & \square \\ \hline \square & \times & \square \\ \hline \end{array} \right) \\
&= \text{ind} \left( - \begin{array}{|c|c|c|c|c|} \hline \square & \square & \square & \square & \times \\ \hline \square & \square & \square & \square & \square \\ \hline \end{array} + (q-1) \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \times & \square \\ \hline \end{array} + (q-1) \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \square & \times & \square & \square \\ \hline \end{array} \right. \\
&\quad \left. - \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \times \\ \hline \square & \square & \square & \square \\ \hline \end{array} + (q^2-1) \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \times \\ \hline \end{array} - q^2 \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \times \\ \hline \end{array} + (q^3-q^2) \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \times \\ \hline \end{array} \right) \\
&= -(q-1) \begin{array}{|c|c|c|c|c|} \hline \square & \square & \square & \square & \square \\ \hline \square & \square & \square & \square & \square \\ \hline \end{array} + (q-1)(q^2-q-1) \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \end{array} \\
&\quad + (q-1)(q^3-q^2) \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \end{array} + (q^2-1)(q^3-q^2) \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \end{array} \\
&\quad - q^2(q^3-q) \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \end{array} + (q^3-q^2)(q^4-q^2) \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \square & \square & \square & \square \\ \hline \end{array}.
\end{aligned}$$

**2.12. Example.** We apply Theorem 2.10 to rederive the explicit formula (1.1) for the Gelfand-Graev character  $\Gamma_n$  itself. Of course, by Theorem 2.10,  $\Gamma_n = \text{ind } \circ e_{n-1}(\pi_{(0)})$ . We will in fact prove that

$$e_{n-1}\pi_{(0)} = (-1)^{n-1} \sum_{\lambda \vdash n} \sum_{B \text{ removable}} (1-q)(1-q^2) \dots (1-q^{h(\lambda)-1}) \pi_{\lambda, B} \quad (2.13)$$

Then (1.1) follows easily on applying  $\text{ind}$  using Lemma 2.9(iv) and the calculation in the proof of Lemma 2.9(v).

To prove (2.13), use induction on  $n$ ,  $n=1$  being immediate from Lemma 2.9(i). For  $n > 1$ , fix some  $\lambda \vdash n$ , label the addable and removable nodes of  $\lambda$  as in Lemma 2.9 and take  $1 \leq r \leq s$ . Thanks to Lemma 2.9(ii), the  $\pi_{\lambda, A_r}$ -coefficient of  $e_n(\pi_{(0)}) = e_+ \circ e_{n-1}(\pi_{(0)})$  only depends on the  $\pi_{\lambda, B_i}$ -coefficients of  $e_{n-1}(\pi_{(0)})$  for  $1 \leq i \leq \min(r, s-1)$ . So by the induction hypothesis the  $\pi_{\lambda, A_r}$ -coefficient of  $e_n(\pi_{(0)})$  is the same as the  $\pi_{\lambda, A_r}$ -coefficient of

$$(-1)^{n-1} (1-q) \dots (1-q^{h(\lambda)-1}) \sum_{i=1}^{\min(r, s-1)} e_+ \pi_{\lambda, B_i},$$

which using Lemma 2.9(ii) equals

$$(-1)^{n-1} (1-q) \dots (1-q^{h(\lambda)-1}) \sum_{i=1}^{\min(r, s-1)} (\delta_{r,i} q^{d(B_i)} - q^{e(B_i)}).$$



This simplifies to  $(-1)^n(1-q)\dots(1-q^{h(\lambda)-1})$  if  $r < s$  and  $(-1)^n(1-q)\dots(1-q^{h(\lambda)})$  if  $r = s$ , as required to prove the induction step.

### 3 The forgotten basis

Recall from the introduction that for  $\lambda \vdash n$ ,  $\chi_\lambda \in C(G_n)$  denotes the irreducible unipotent character parametrized by the partition  $\lambda$ , and  $\sigma_\lambda \in \mathfrak{g}$  is its projection to unipotent-supported class functions. Also,  $\{\vartheta_\lambda\}$  denotes the ‘forgotten’ basis of  $\mathfrak{g}$ , which can be defined as the unique basis of  $\mathfrak{g}$  such that for each  $n$  and each  $\lambda \vdash n$ ,

$$\sigma_{\lambda'} = \sum_{\mu \vdash n} K_{\lambda, \mu} \vartheta_\mu. \quad (3.1)$$

Given  $\lambda \vdash n$ , we write  $\mu \perp_j \lambda$  if  $\mu \vdash (n-j)$  and  $\lambda_{i+1} \leq \mu_i \leq \lambda_i$  for all  $i = 1, 2, \dots$ . This definition arises in the following well-known inductive formula for the Kostka number  $K_{\lambda, \mu}$ , i.e. the number of standard  $\lambda$ -tableaux of weight  $\mu$  [9, §I(6.4)]:

**3.2. Lemma.** *For  $\lambda \vdash n$  and any composition  $\nu \vDash n$ ,  $K_{\lambda, \nu} = \sum_{\mu \perp_j \lambda} K_{\mu, \bar{\nu}}$ , where  $j$  is the last non-zero part of  $\nu$  and  $\bar{\nu}$  is the composition obtained from  $\nu$  by replacing this last non-zero part by zero.*

We will need the following special case of Zelevinsky’s branching rule [13, §13.5] (see also [1, Corollary 5.4d(ii)] for its modular analogue):

**3.3. Theorem (Zelevinsky).** *For  $\lambda \vdash n$ ,  $\text{res}_{H_{n-1}}^{G_n} \chi_{\lambda'} = \sum_{j \geq 1} \sum_{\mu \perp_j \lambda} e_{j-1} \chi_{\mu'}$ .*

Now define the map  $\delta : \mathfrak{g} \rightarrow \mathfrak{g}$  as in the introduction by setting  $\delta(\pi_\lambda) = \frac{1}{q^{h(\lambda)} - 1} \pi_\lambda$  for each partition  $\lambda$  and extending linearly to all of  $\mathfrak{g}$ . The significance of  $\delta$  is that by Lemma 2.9(v),  $\delta \circ \text{ind} \circ \text{res}(\pi_\lambda) = \pi_\lambda$  for all  $\lambda$ . Set  $\hat{\gamma}_n = \delta \circ \gamma_n$  and for a partition  $\lambda$ , define

$$\hat{\gamma}_\lambda = \sum_{(n_1, \dots, n_h)} \hat{\gamma}_{n_h} \circ \dots \circ \hat{\gamma}_{n_2} \circ \hat{\gamma}_{n_1} \quad (3.4)$$

summing over all compositions  $(n_1, \dots, n_h)$  obtained by reordering the  $h = h(\lambda)$  non-zero parts of  $\lambda$  in all possible ways.

**3.5. Theorem.** *For any  $\lambda \vdash n$ ,  $\vartheta_\lambda = \hat{\gamma}_\lambda(\pi_{(0)})$ .*

*Proof.* We will show by induction on  $n$  that

$$\sigma_{\lambda'} = \sum_{\mu \vdash n} K_{\lambda, \mu} \hat{\gamma}_\mu(\pi_{(0)}). \quad (3.6)$$

The theorem then follows immediately in view of the definition (3.1) of  $\vartheta_\lambda$ . Our induction starts trivially with the case  $n = 0$ . So now suppose that  $n > 0$  and that (3.6) holds for all smaller  $n$ . By Theorem 3.3,

$$\text{res } \sigma_{\lambda'} = \sum_{j \geq 1} \sum_{\mu \perp_j \lambda} e_{j-1} \sigma_{\mu'}.$$

Applying the operator  $\delta \circ \text{ind}$  to both sides, we deduce that

$$\sigma_{\lambda'} = \sum_{j \geq 1} \sum_{\mu \perp_j \lambda} \hat{\gamma}_j(\sigma_{\mu'}) = \sum_{j \geq 1} \sum_{\mu \perp_j \lambda} \sum_{\nu \vdash (n-j)} K_{\mu, \nu} \hat{\gamma}_j \circ \hat{\gamma}_\nu(\pi_{(0)})$$

(we have applied the induction hypothesis)

$$= \sum_{j \geq 1} \sum_{\mu \perp_j \lambda} \sum_{\nu \vdash (n-j)} \sum_{(n_1, \dots, n_h)} K_{\mu, (n_1, \dots, n_h)} \hat{\gamma}_j \circ \hat{\gamma}_{n_h} \circ \dots \circ \hat{\gamma}_{n_2} \circ \hat{\gamma}_{n_1}(\pi_{(0)})$$

(summing over  $(n_1, \dots, n_h)$  obtained by reordering the non-zero parts  $\nu$  in all possible ways)

$$= \sum_{j \geq 1} \sum_{\nu \vdash (n-j)} \sum_{(n_1, \dots, n_h)} K_{\lambda, (n_1, \dots, n_h, j)} \hat{\gamma}_j \circ \hat{\gamma}_{n_h} \circ \dots \circ \hat{\gamma}_{n_2} \circ \hat{\gamma}_{n_1}(\pi_{(0)})$$

(we have applied Lemma 3.2)

$$= \sum_{\eta \vdash n} \sum_{(m_1, \dots, m_k)} K_{\lambda, (m_1, \dots, m_k)} \hat{\gamma}_{m_k} \circ \dots \circ \hat{\gamma}_{m_2} \circ \hat{\gamma}_{m_1}(\pi_{(0)})$$

(now summing over  $(m_1, \dots, m_k)$  obtained by reordering the non-zero parts of  $\eta$  in all possible ways)

$$= \sum_{\eta \vdash n} K_{\lambda, \eta} \hat{\gamma}_\eta(\pi_{(0)})$$

which completes the proof.  $\square$

**3.7. Example.** For  $\chi \in \mathfrak{g}_n$ , write  $\deg \chi$  for its value at the identity element of  $G_n$ . We wish to derive the formula (1.7) for  $\deg \vartheta_\lambda = S_{\lambda, (1^n)}(q)$  using Theorem 3.5. So, fix  $\lambda \vdash n$ . Then, by Theorem 3.5,

$$\deg \vartheta_\lambda = \sum_{(n_1, \dots, n_h)} \deg [\hat{\gamma}_{n_h} \circ \dots \circ \hat{\gamma}_{n_1}(\pi_{(0)})]. \quad (3.8)$$

We will show that given  $\chi \in \mathfrak{g}_m$ ,

$$\deg \hat{\gamma}_n(\chi) = (q^{m+n-1} - 1)(q^{m+n-2} - 1) \dots (q^{m+1} - 1) \deg \chi; \quad (3.9)$$

then the formula (1.7) follows easily from (3.8). Now  $\gamma_n$  is Harish-Chandra multiplication by  $\Gamma_n$ , so

$$\deg \gamma_n(\chi) = \deg \Gamma_n \deg \chi \cdot \frac{(q^{m+n} - 1) \dots (q^{m+1} - 1)}{(q^n - 1) \dots (q - 1)},$$

the last term being the index in  $G_{m+n}$  of the standard parabolic subgroup with Levi factor  $G_m \times G_n$ . This simplifies using (1.1) to  $(q^{m+n} - 1) \dots (q^{m+1} - 1) \deg \chi$ . Finally, to calculate  $\deg \hat{\gamma}_n(\chi)$ , we need to rescale using  $\delta$ , which divides this expression by  $(q^{m+n} - 1)$ .

## 4 Brauer character values

Finally, we derive the formula (1.5) for the unipotent Brauer character values. So let  $p$  be a prime not dividing  $q$  and  $\chi_\lambda^p$  be the irreducible unipotent  $p$ -modular Brauer character labelled by  $\lambda$  as in the introduction. Writing  $C^p(G_n)$  for the  $\mathbb{C}$ -valued class functions on  $G_n$  with support in the set of  $p'$ -elements of  $G_n$ , we view  $\chi_\lambda^p$  as an element of  $C^p(G_n)$ . Let  $\dot{\chi}_\lambda$  denote the projection of the ordinary unipotent character  $\chi_\lambda$  to  $C^p(G_n)$ . Then, by [6], we can write

$$\dot{\chi}_\lambda = \sum_{\mu \vdash n} D_{\lambda, \mu} \chi_\mu^p, \quad (4.1)$$

and the resulting matrix  $D = (D_{\lambda, \mu})$  is the *unipotent part* of the  $p$ -modular decomposition matrix of  $G_n$ . One of the main achievements of the Dipper-James theory from [3] (see e.g. [1, (3.5a)]) relates these decomposition numbers to the decomposition numbers of quantum  $GL_n$ .

To recall some definitions, let  $\mathbb{k}$  be a field of characteristic  $p$  and  $v \in \mathbb{k}$  be a square root of the image of  $q$  in  $\mathbb{k}$ . Let  $U_n$  denote the divided power version of the quantized enveloping algebra  $U_q(\mathfrak{gl}_n)$  specialized over  $\mathbb{k}$  at the parameter  $v$ , as defined originally by Lusztig [8] and Du [4, §2] (who extended Lusztig's construction from  $\mathfrak{sl}_n$  to  $\mathfrak{gl}_n$ ). For each partition  $\lambda \vdash n$ , there is an associated irreducible polynomial representation of  $U_n$  of high-weight  $\lambda$ , which we denote by  $L(\lambda)$ . Also let  $V(\lambda)$  denote the standard (or Weyl) module of high-weight  $\lambda$ . Write

$$\text{ch } V(\lambda) = \sum_{\mu \vdash n} D'_{\lambda, \mu} \text{ch } L(\mu), \quad (4.2)$$

so  $D' = (D'_{\lambda, \mu})$  is the decomposition matrix for the polynomial representations of quantum  $GL_n$  of degree  $n$ . Then, by [3]:

**4.3. Theorem (Dipper and James).**  $D'_{\lambda, \mu} = D_{\lambda', \mu'}$ .

Let  $\sigma_\lambda^p$  denote the projection of  $\chi_\lambda^p$  to unipotent-supported class functions. The  $\{\sigma_\lambda^p\}$  also give a basis for the Hall algebra  $\mathfrak{g}$ . Inverting (4.1) and using (3.1),

$$\sigma_{\lambda'}^p = \sum_{\mu \vdash n} D_{\lambda', \mu'}^{-1} \sigma_{\mu'} = \sum_{\mu, \nu \vdash n} D_{\lambda', \mu'}^{-1} K_{\mu, \nu} \vartheta_\nu \quad (4.4)$$

where  $D^{-1} = (D_{\lambda,\mu}^{-1})$  is the inverse of the matrix  $D$ . On the other hand, writing  $K_{\lambda,\mu}^{p,\ell}$  for the multiplicity of the  $\mu$ -weight space of  $L(\lambda)$ , and recalling that  $K_{\lambda,\mu}$  is the multiplicity of the  $\mu$ -weight space of  $V(\lambda)$ , we have by (4.2) that

$$K_{\mu,\nu} = \sum_{\eta \vdash n} D'_{\mu,\eta} K_{\eta,\nu}^{p,\ell}. \quad (4.5)$$

Substituting (4.5) into (4.4) and applying Theorem 4.3, we deduce:

4.6. **Theorem.**  $\sigma_{\lambda'}^p = \sum_{\mu \vdash n} K_{\lambda,\mu}^{p,\ell} \vartheta_{\mu}.$

Now (1.5) and (1.6) follow at once. This completes the proof of the formulae stated in the introduction.

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